

# Fourierova optika

Fourierova transformace

$$f(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} g(\omega) \exp(-i\omega t) d\omega,$$

$$g(\omega) = \int_{-\infty}^{\infty} f(t) \exp(+i\omega t) dt$$

$$f(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} g(k) \exp(ikx) dk \quad g(k) = \int_{-\infty}^{\infty} f(x) \exp(-ikx) dx$$

Rozvoj do rovinných vln

$$f(x, y) = \frac{1}{(2\pi)^2} \int_{-\infty}^{\infty} g(k_x, k_y) \exp[i(k_x x + k_y y)] dk_x dk_y$$

$$g(k_x, k_y) = \int_{-\infty}^{\infty} f(x, y) \exp[-i(k_x x + k_y y)] dx dy$$

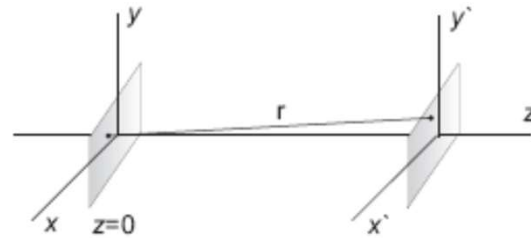
$$k_x = \frac{2\pi}{\lambda_x}, \quad k_y = \frac{2\pi}{\lambda_y}$$

Prostorové frekvence

Prostorové periody

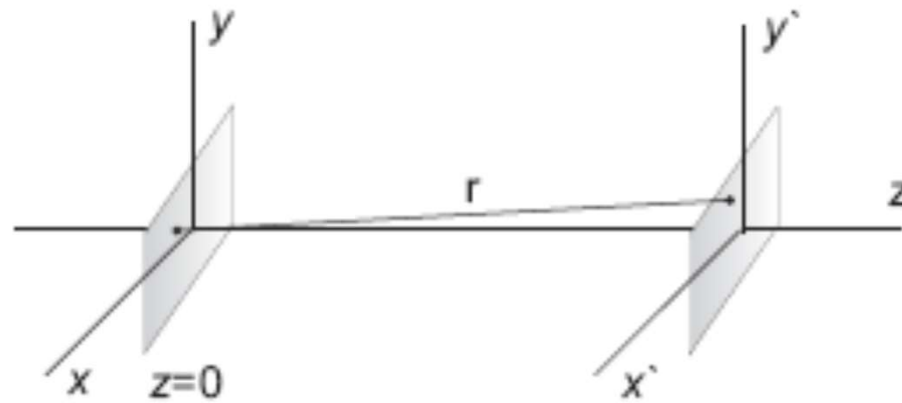
## Fresnelova a Fraunhoferova aproximace

Znamé pole E v  $z=0$



$$\tilde{E}(x', y') = \frac{-i}{\lambda} A \int_{\text{APERTURA}} F(\theta) \frac{1}{r r_z} \exp[i k (r + r_z)] dS$$

$$\tilde{E}(x', y') = \frac{-i}{\lambda} \int_{\text{APERTURA}} \tilde{E}_S(x, y) \frac{1}{r} \exp(i k r) dx dy$$



$$r = \sqrt{z^2 + (x-x')^2 + (y-y')^2} \approx z + \frac{(x-x')^2 + (y-y')^2}{2z}.$$

PARAXIÁLNÍ APROXIMACE:

$F=1, r \rightarrow z$  ve jmenovateli

$$\tilde{E}(x', y') =$$

$$= \frac{-i}{\lambda} \frac{1}{z} \exp(ikz) \int_{\text{APERTURA}} \tilde{E}_s(x, y) \exp\left\{\frac{ik}{2z}[(x-x')^2 + (y-y')^2]\right\} dx dy,$$

$$\tilde{E}(x', y') = \frac{-i}{\lambda} \frac{1}{z} \exp(ikz) \exp\left[\frac{ik}{2z}(x'^2 + y'^2)\right]$$

$$\int_{\text{APERTURA}} \tilde{E}_s(x, y) \exp\left[\frac{ik}{2z}(x^2 + y^2)\right] \exp\left[\frac{-ik}{z}(xx' + yy')\right] dx dy.$$

**Fresnelova paraxiální aproximace**

$$\begin{aligned}
\tilde{E}(x', y') &= \\
&= \frac{-i}{\lambda} \frac{1}{z} \exp(ikz) \int_{\text{APERTURA}} \tilde{E}_s(x, y) \exp\left\{\frac{ik}{2z}[(x-x')^2 + (y-y')^2]\right\} dx dy, \\
\tilde{E}(x', y') &= \frac{-i}{\lambda} \frac{1}{z} \exp(ikz) \exp\left[\frac{ik}{2z}(x'^2 + y'^2)\right] \\
&\quad \int_{\text{APERTURA}} \tilde{E}_s(x, y) \exp\left[\frac{ik}{2z}(x^2 + y^2)\right] \exp\left[\frac{-ik}{z}(xx' + yy')\right] dx dy. \\
&\quad = 0
\end{aligned}$$

**Fresnelova paraxiální aproximace**

**Fraunhoferova aproximace**

$$\tilde{E}(x', y') =$$

$$= \frac{-i}{\lambda} \frac{1}{z} \exp(ikz) \exp\left[\frac{ik}{2z}(x'^2 + y'^2)\right] \int_{\text{APERATURA}} \tilde{E}_S(x, y) \exp\left[\frac{-ik}{z}(xx' + yy')\right] dx dy$$

$$k_x = \frac{kx'}{z}, \quad k_y = \frac{ky'}{z}.$$

Zklenutí pole, pro body blízko osy možno brát konst.

$$\tilde{E}(x', y') = \tilde{E}(k_x, k_y) = \tilde{A} \int_{\text{APERATURA}} \tilde{E}_S(x, y) \exp[-i(k_x x + k_y y)] dx dy$$

**Prostorové frekvence**  $k_x = \frac{kx'}{z}, \quad k_y = \frac{ky'}{z}.$



$$\tilde{E}(x', y') =$$

$$= \frac{-i}{\lambda} \frac{1}{z} \exp(ikz) \exp\left[\frac{ik}{2z}(x'^2 + y'^2)\right] \int_{\text{APERATURA}} \tilde{E}_S(x, y) \exp\left[\frac{-ik}{z}(xx' + yy')\right] dx dy$$

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Zklenutí pole, pro body blízko osy možno brát konst.

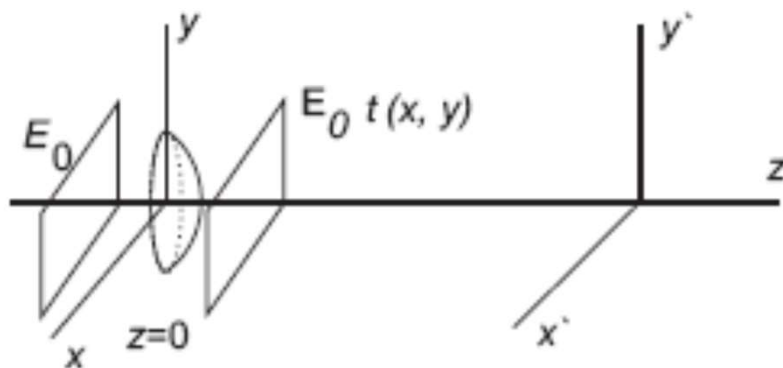
$$\tilde{E}(x', y') = \tilde{E}(k_x, k_y) = \tilde{A} \int_{\text{APERATURA}} \tilde{E}_S(x, y) \exp[-i(k_x x + k_y y)] dx dy$$

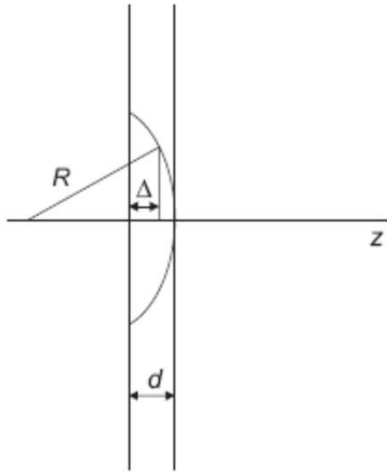
## Fourierova transformace

Prostorové frekvence  $k_x = \frac{kx'}{z}, \quad k_y = \frac{ky'}{z}.$

Daleké pole (abychom mohli zanedbat ten člen s kvadráty  $x'$ ,  $y'$ ) nutné pro Fraunhoferovu aproximaci.

Co teď **čočka??**





$$\tilde{E}_s(x, y) = \tilde{E}_0 \tilde{t}(x, y)$$

$$\tilde{t}(x, y) = \exp[i\varphi(x, y)]$$

$$\varphi(x, y) = k_0 d + (n-1)k_0 \Delta(x, y)$$

$$\Delta(x, y) = d - \left( R - \sqrt{R^2 - x^2 - y^2} \right)$$

$$\Delta(x, y) = d - R \left( 1 - \sqrt{1 - \frac{x^2 + y^2}{R^2}} \right)$$

$$x^2 + y^2 \ll R^2$$

$$\sqrt{1 - \frac{x^2 + y^2}{R^2}} \approx 1 - \frac{x^2 + y^2}{2R^2}$$

$$\Delta(x, y) = d - \frac{x^2 + y^2}{2R}$$

$$\tilde{t}(x, y) = \exp(i\varphi_0) \exp\left(-ik_0 \frac{x^2 + y^2}{2f'}\right)$$

kde

$$\frac{1}{f'} = \frac{n-1}{R}$$

$$\begin{aligned} \tilde{E}(x', y') &= \\ &= \frac{-i}{\lambda z} \exp(i k z + i \varphi_0) \exp\left(ik \frac{x'^2 + y'^2}{2z}\right) \\ &\iint_{APERTURA} \tilde{E}_0 \exp\left[ik \frac{x^2 + y^2}{2} \left(\frac{1}{z} - \frac{1}{f'}\right)\right] \exp\left[-\frac{ik}{z}(x x' + y y')\right] dx dy \end{aligned}$$

Osvětlení čočky může být modulované  $\tilde{E}_0 = \tilde{E}_0(x, y)$

V ohniskové rovině čočky  $z = f'$

$$\tilde{E}(x', y') = \tilde{E}(k_x, k_y) = \tilde{B} \int_{APERTURA} \tilde{E}_S(x, y) \exp[-i(k_x x + k_y y)] dx dy$$

$$\begin{aligned} \tilde{E}(x', y') &= \\ &= \frac{-i}{\lambda z} \exp(i k z + i \varphi_0) \exp\left(ik \frac{x'^2 + y'^2}{2 z}\right) \\ &\iint_{APERTURA} \tilde{E}_0 \exp\left[ik \frac{x^2 + y^2}{2} \left(\frac{1}{z} - \frac{1}{f'}\right)\right] \exp\left[-\frac{ik}{z}(x x' + y y')\right] dx dy \end{aligned}$$

vynuluje se

Osvětlení čočky může být modulované

$$\tilde{E}_0 = \tilde{E}_0(x, y)$$

V ohniskové rovině čočky

$$z = f'$$

$$\tilde{E}(x', y') = \tilde{E}(k_x, k_y) = \tilde{B} \int_{APERTURA} \tilde{E}_S(x, y) \exp[-i(k_x x + k_y y)] dx dy$$

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## Opět Fourierova transformace !

Ted' bez ohledu na daleké pole !!!

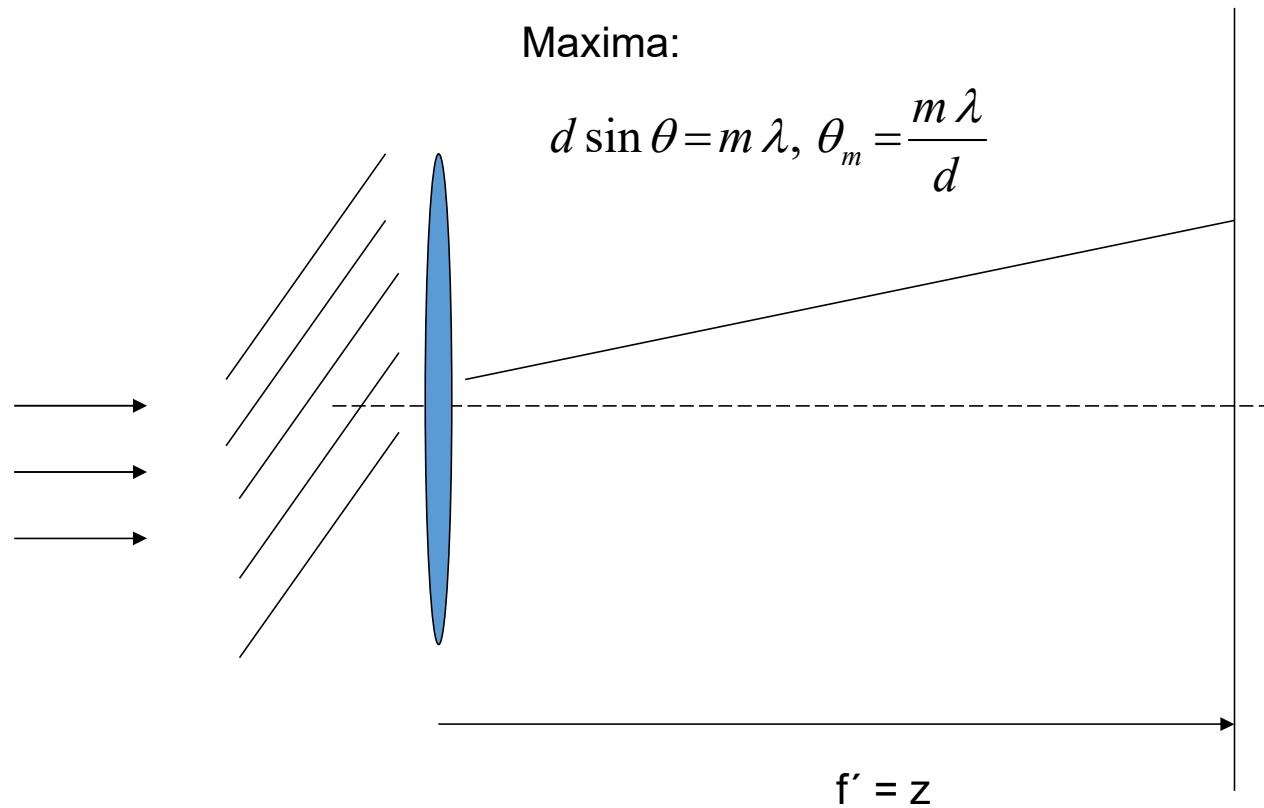
Čili pro ohniskové roviny čoček možno používat Fraunhoferovu difrakci !!!!!

Například: teleskop, mikroskop, oko, apod.

# Prostorová filtrace

Příklad optického zpracování obrazu





$$\theta_m = y'_m / f'$$

$$y'_m = m \frac{\lambda f'}{d}$$

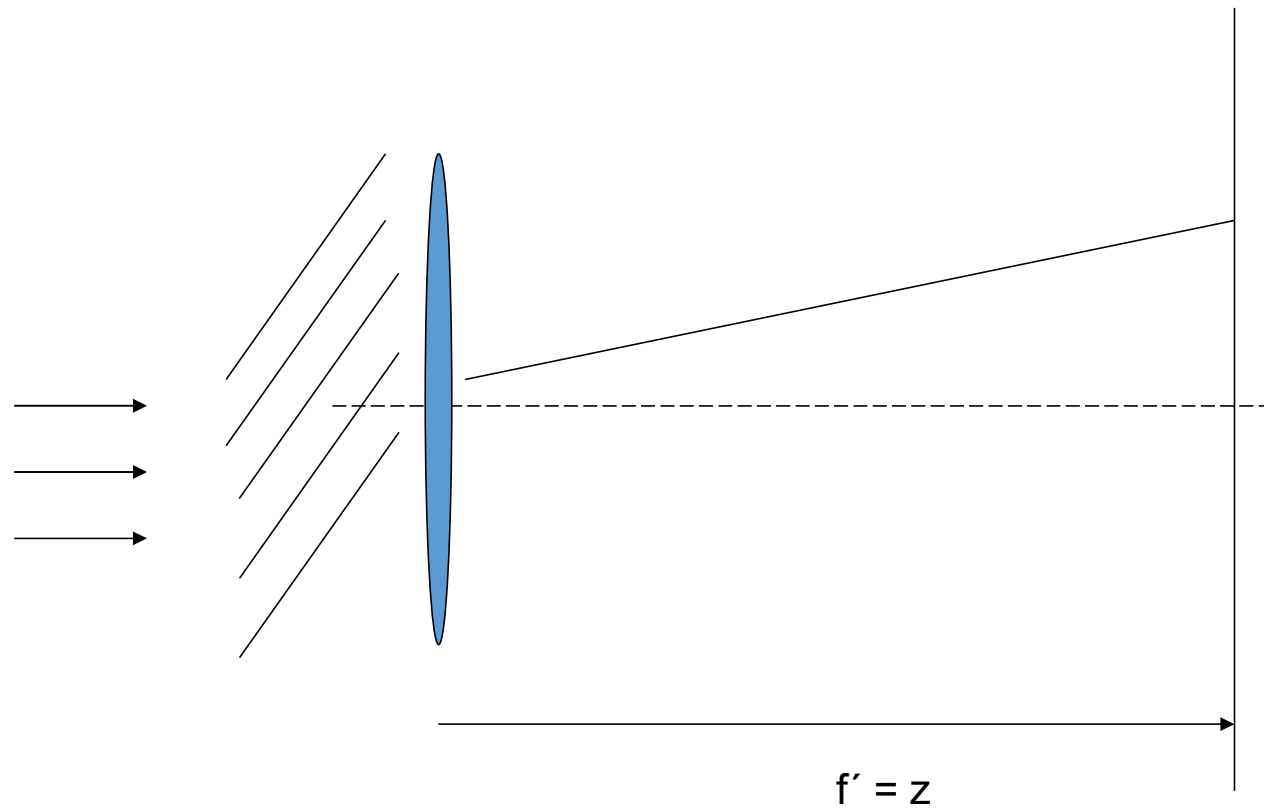
Prostorová frekvence

$$k_y = \frac{k y'}{z}$$

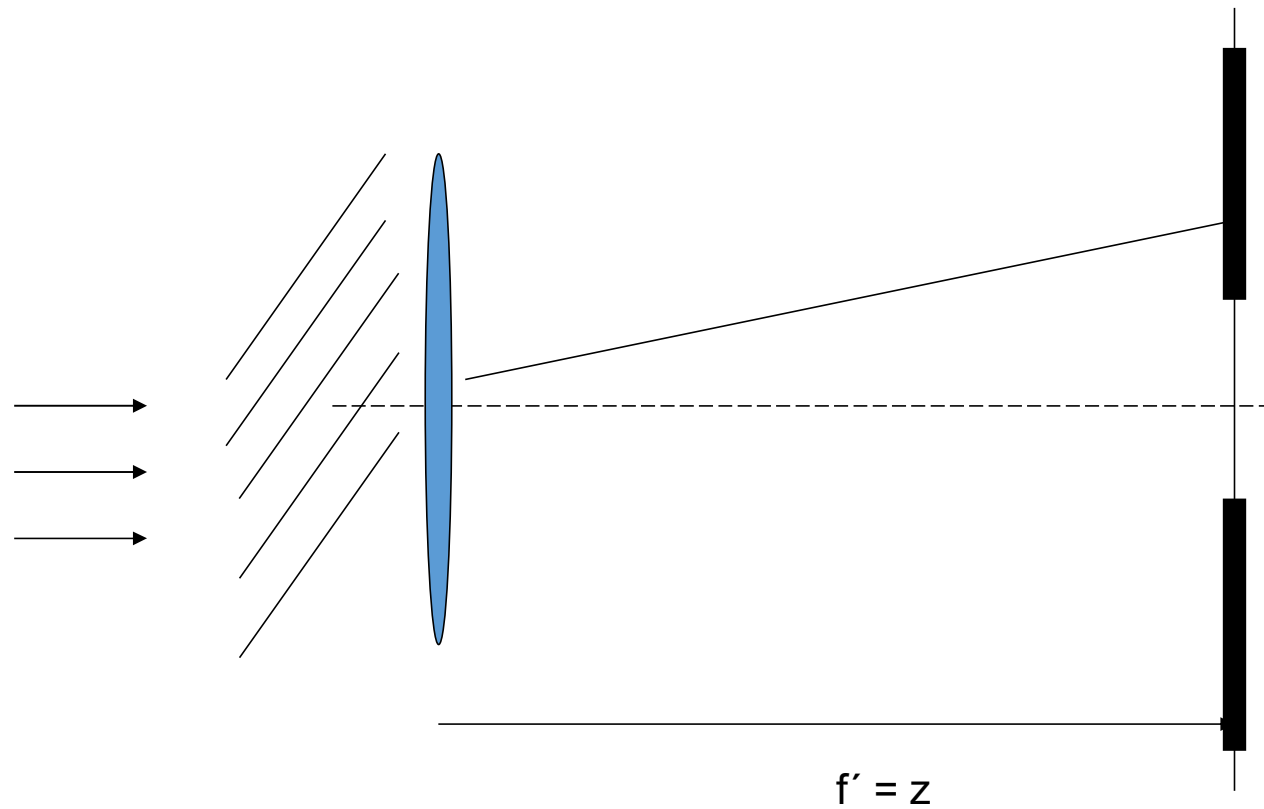
$$k_{ym} = \frac{m 2 \pi}{d}$$

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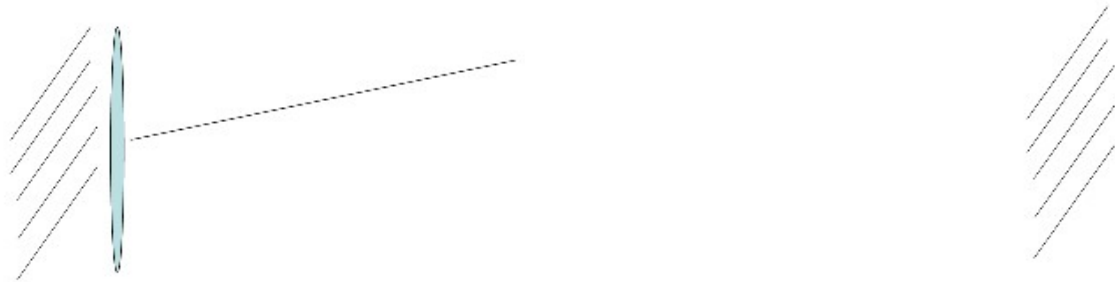
Menší mřížková konstanta  $d$ , větší frekvence (větší úhel)



$$k_{ym} = \frac{m 2 \pi}{d}$$



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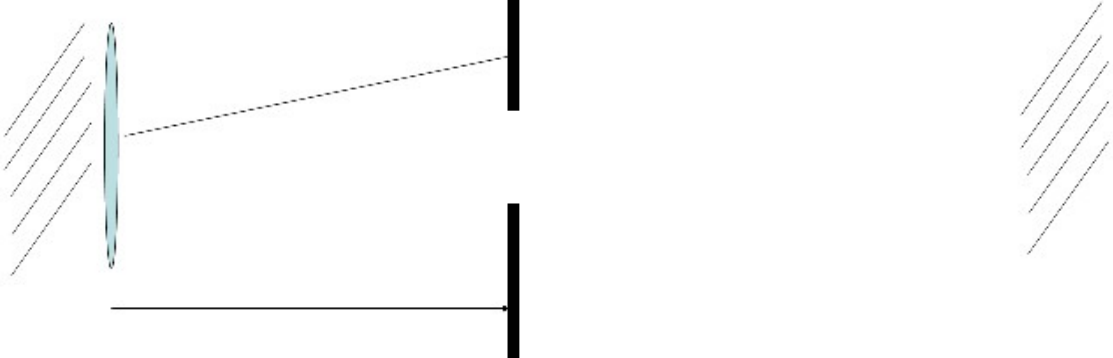


Four. trafo

Zpětná Four. trafo

**clona**

FILTRACE - pustí se jen něco



Four. trafo

Zpětná Four. trafo

# **Optické zpracování obrazu**

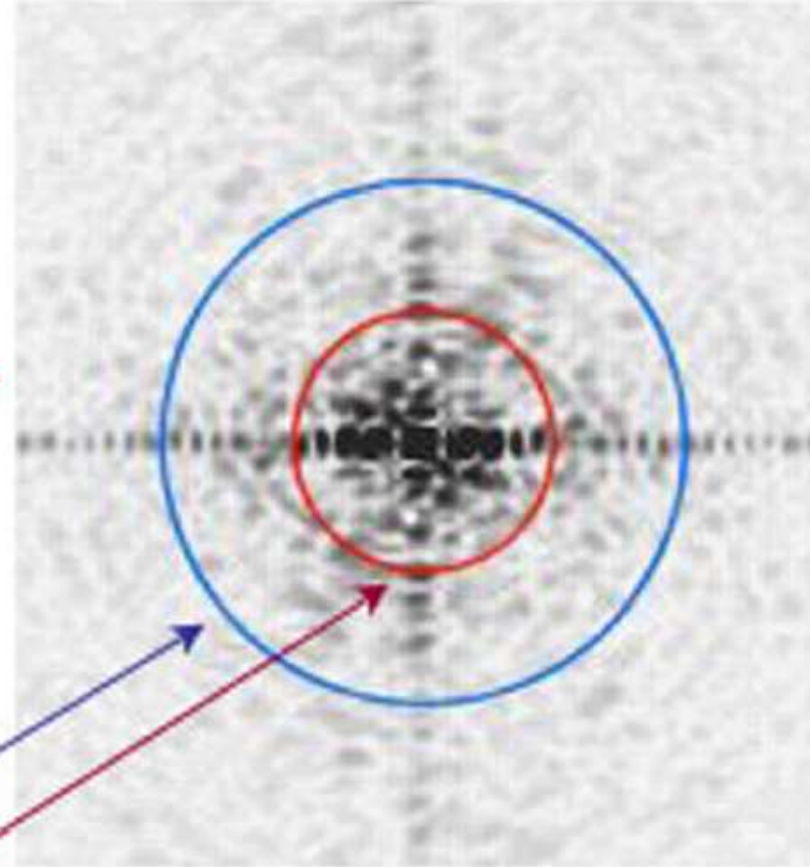
**Experiment: J. Preclíková a K. Žídek**



# František Palacký



FT  
→



$d = 2,12 \text{ mm}$

$d = 1,06 \text{ mm}$





d=2,12mm

d=1,91mm

d=1,69mm

d=1,48mm



d=1,27mm

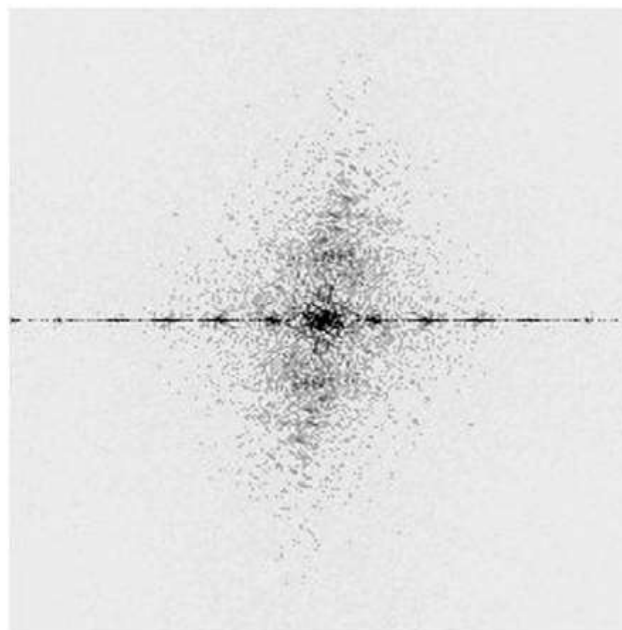
d=1,01mm

d=0,85mm

d=0,42mm

d=0,21mm

## Tygr v kleci



**FT** 



